



Low-temperature quantum transport in CVD-grown single crystal graphene

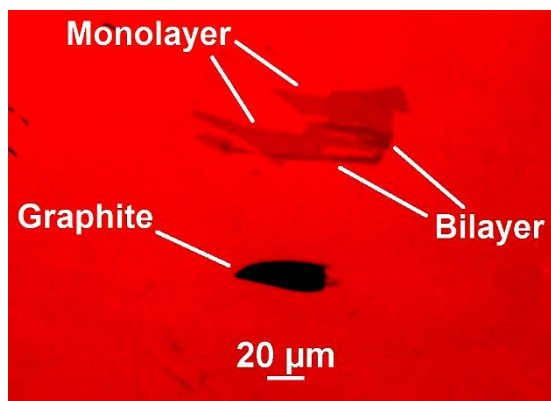
Stefano Guiducci

NanoSEA 2016, Giardini Naxos

Introduction

Exfoliated graphene

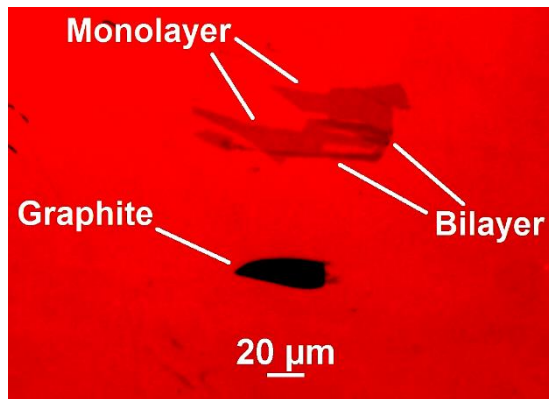
- Excellent quality.
- No need for special equipment.
- Small flakes, random shape.
- Poor scalability.



Introduction

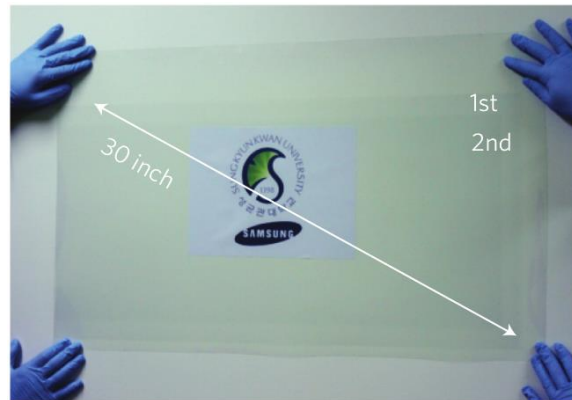
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CVD-grown graphene

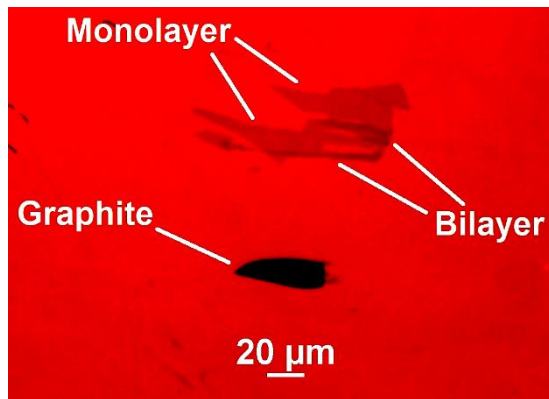
- Excellent size.
- Fast and reliable growth.
- Excellent scalability.
- Lower graphene quality (polycrystalline, defects and contamination).



Introduction

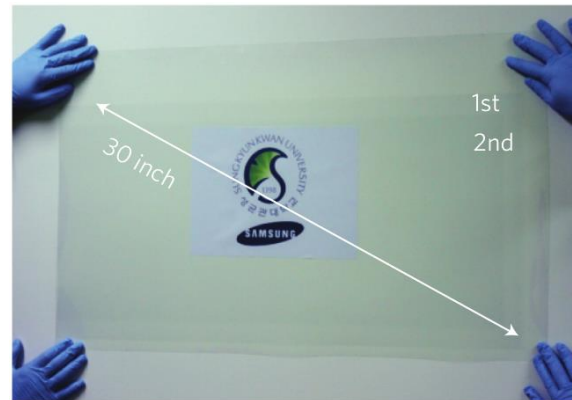
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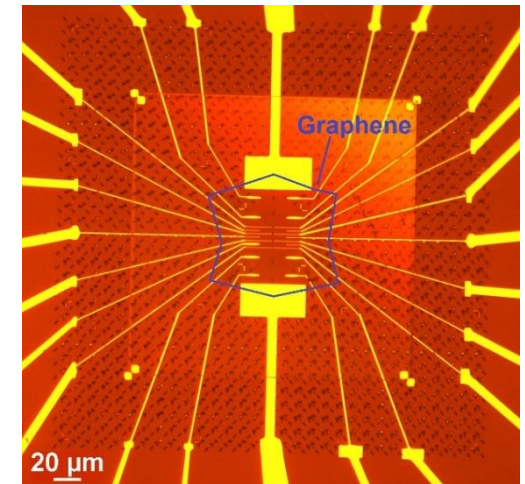
CVD-grown graphene

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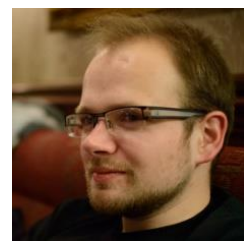
CVD-grown single crystal graphene

- Excellent quality.
- Excellent size and regular shape.
- Fast and reliable growth.
- Excellent scalability.



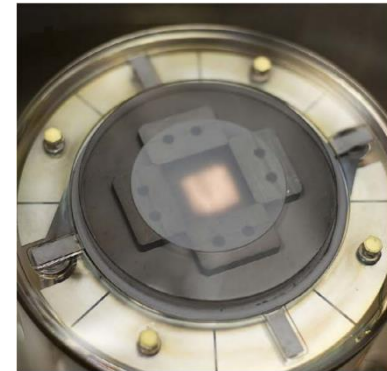
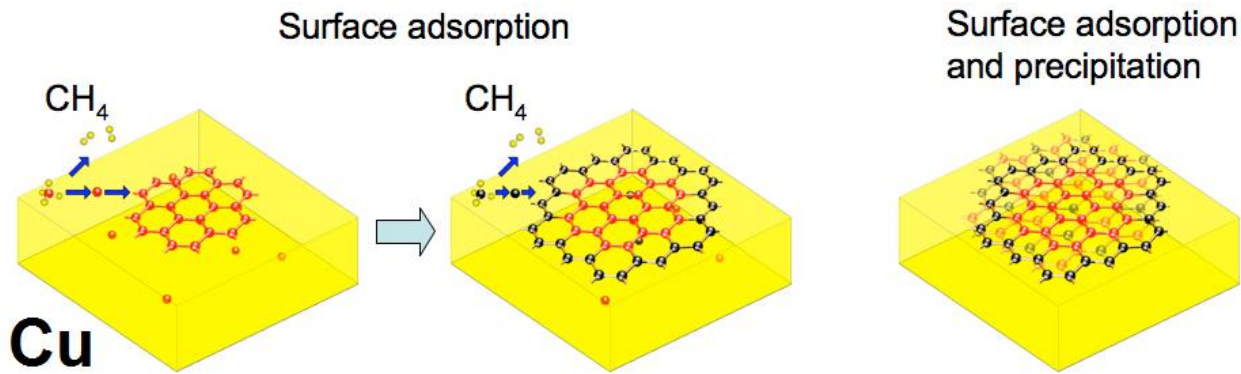
We fabricated several devices that we tested in magneto-transport at low temperature to prove its good electronic quality.

CVD graphene growth

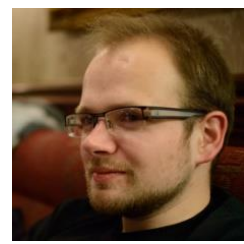


V. Miseikis

- Graphene grown on copper.
- Chemical reaction: $CH_4 \rightarrow 2H_{2(g)} + C_{(s)}$
- Growth in a commercial reactor to increase reproducibility.
- Improvements:
 - Number of nucleation sites drastically reduced (oxidized copper).
 - Semi-dry technique to detach graphene from copper (electrochemical delamination).

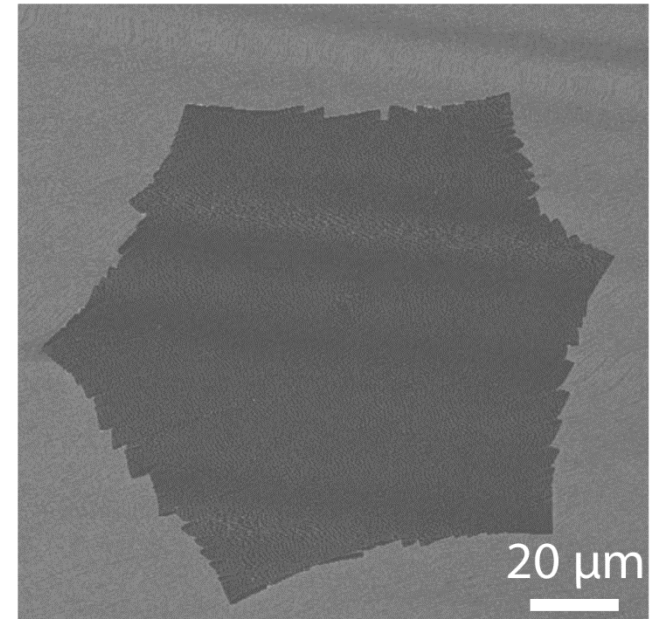
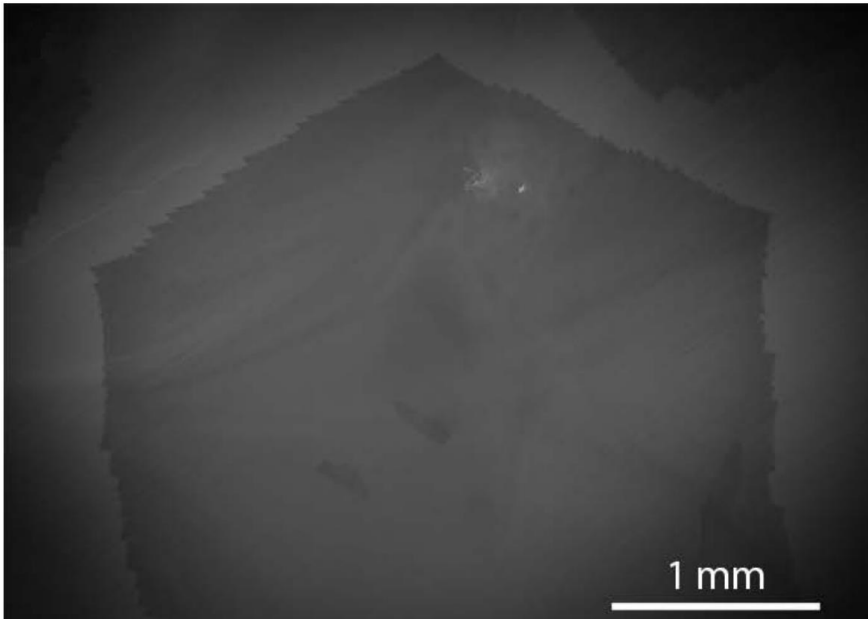


CVD graphene growth



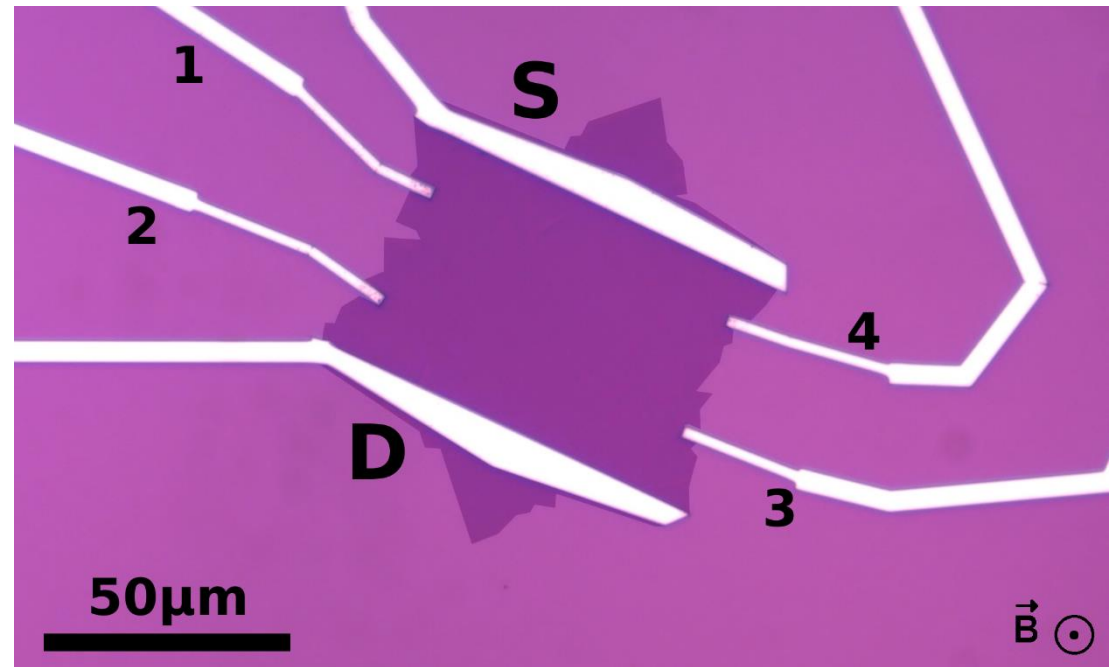
V. Miseikis

- Monocrystalline graphene flakes extended up to several millimeters.
- Fast growth (1mm per hour).
- High crystalline quality already shown with SEM, TEM, Raman, SAED, LEED, XPS.



Electron transport characterization

- Graphene transferred on 300nm of SiO₂ (n++ doped Si substrate used as a backgate).
- Hallbar fabricated (approx. 50μm x70μm).



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- Hallbar fabricated (approx. 50μm x70μm).
- Four voltage probes used to measure longitudinal and transversal (Hall) resistance, respectively:

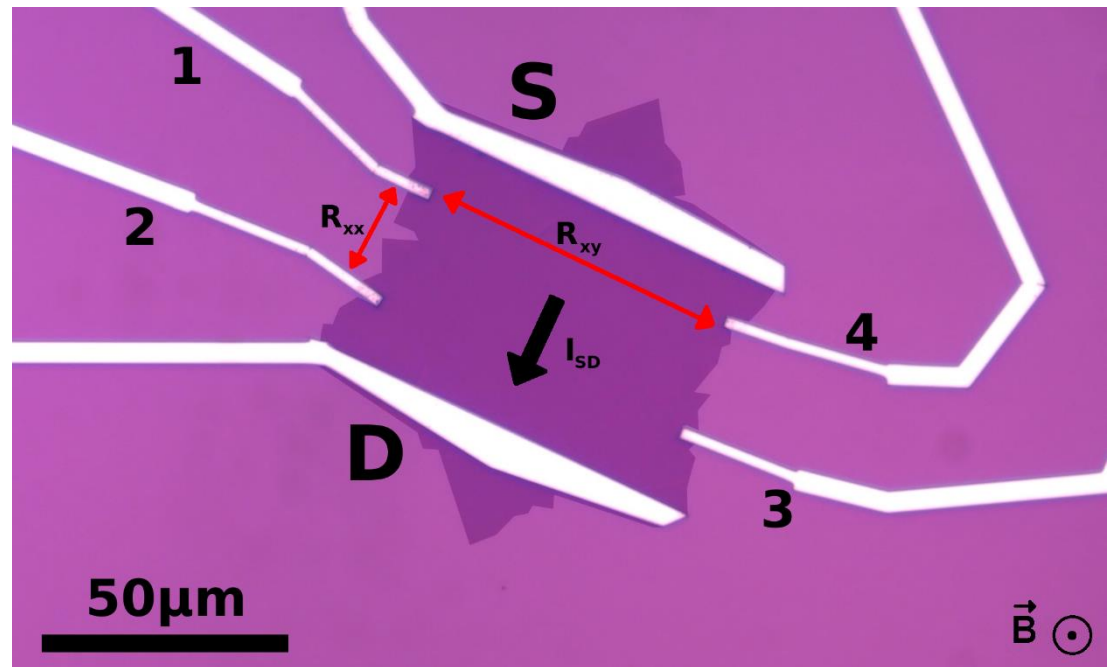
$$R_{xx} = V_{12} / I_{SD}$$

$$R_{xy} = V_{14} / I_{SD}$$

- No sample annealing in order to measure the graphene «as is».
- He-3 cryostat, temperature range:

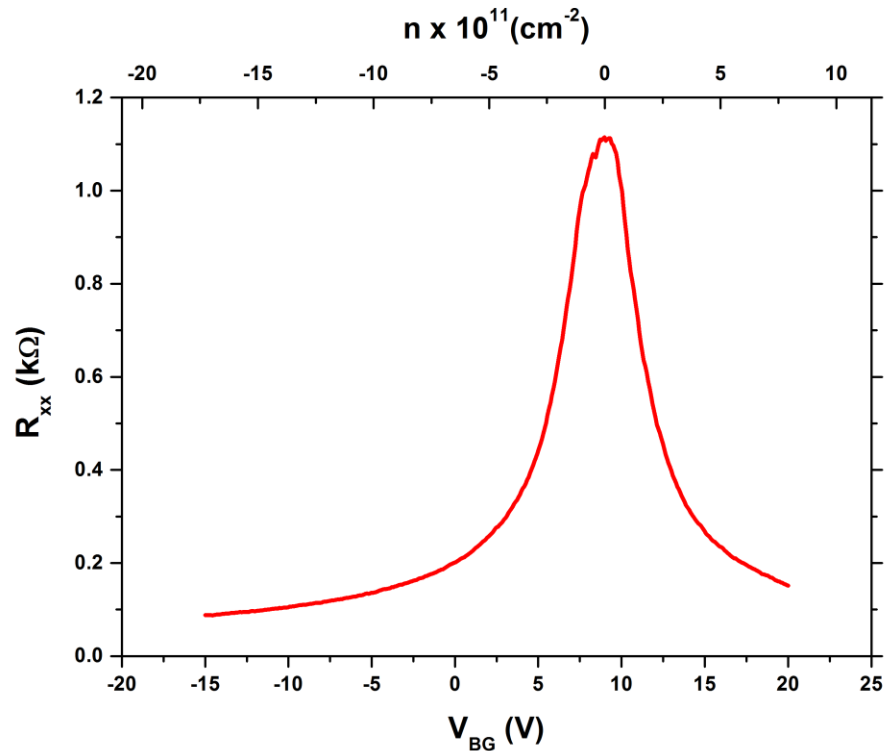
 0.25K - 20K
- Out of plane magnetic field, range:

 0 - 10 Tesla

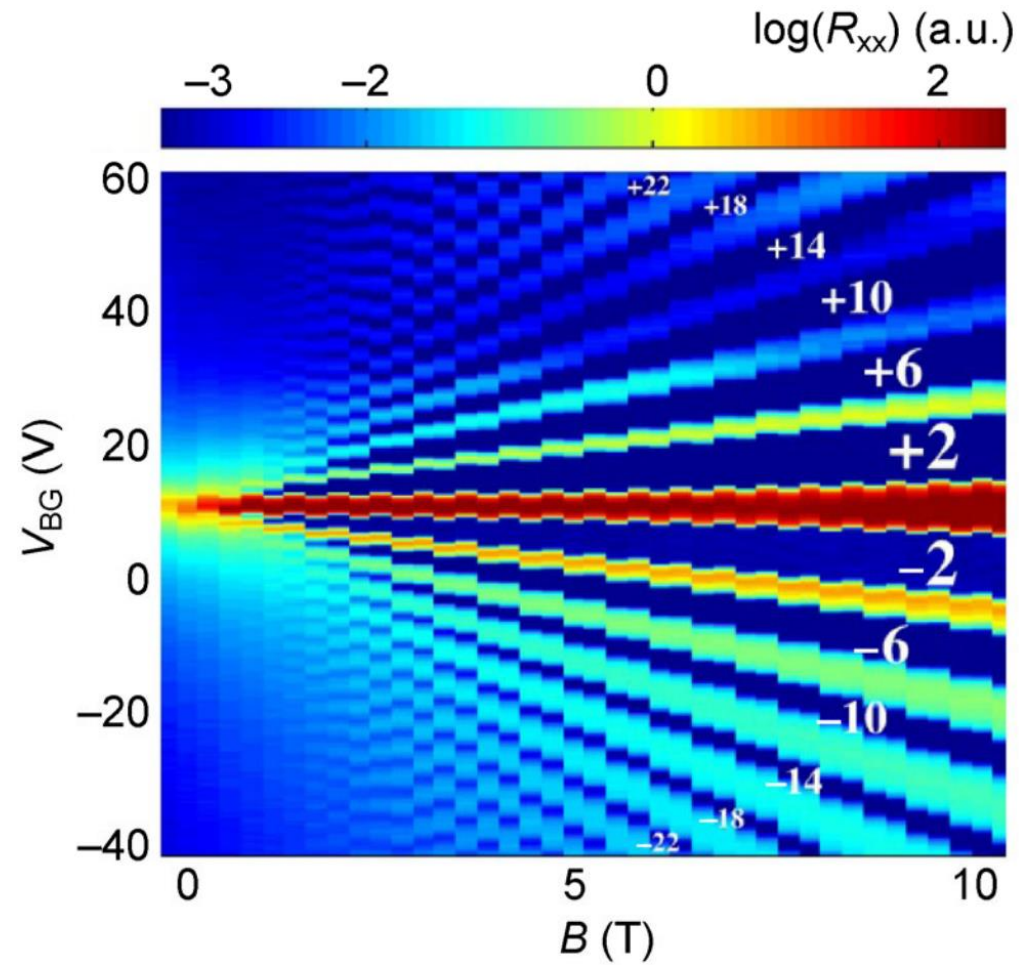
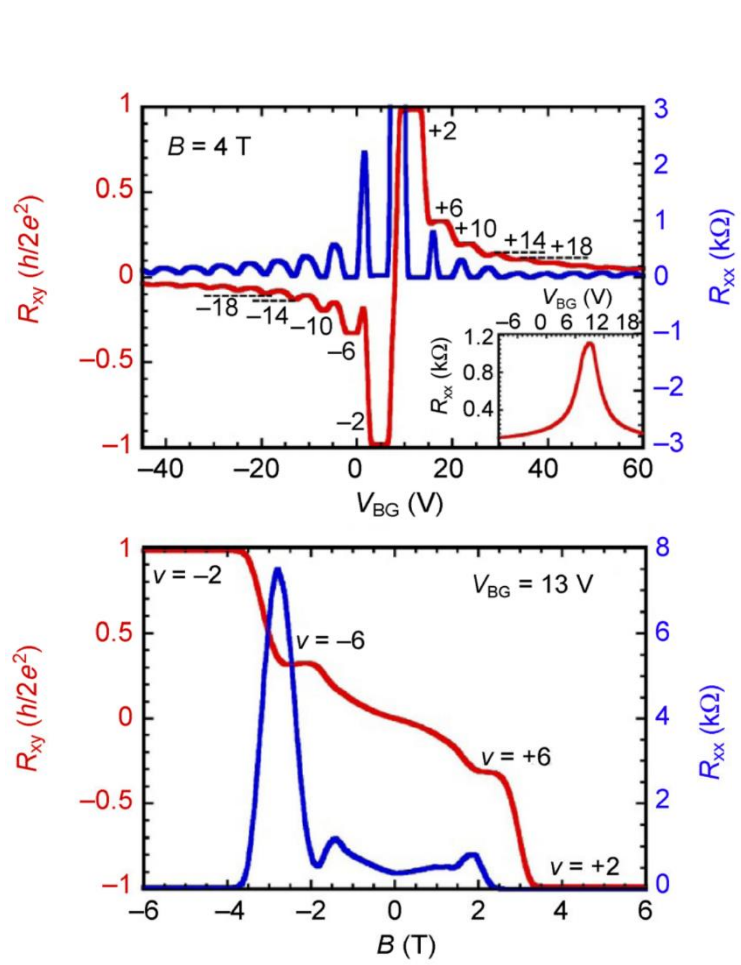


Electron transport characterization

- $\mu = 1.13 \times 10^4 \text{ cm}^2/\text{Vs}$ (mobility at 250mK).
- $V_{\text{Dirac}} = +9\text{V}$
- $n(V_{\text{BG}}=0) = -6.5 \times 10^{11} \text{ cm}^{-2}$ (p-doped) \rightarrow
 \rightarrow Low intrinsic carrier concentration
 (in fact values in the 10^{12} cm^{-2} range are reported for CVD polycrystalline graphene and with the substrate removed by wet etching).



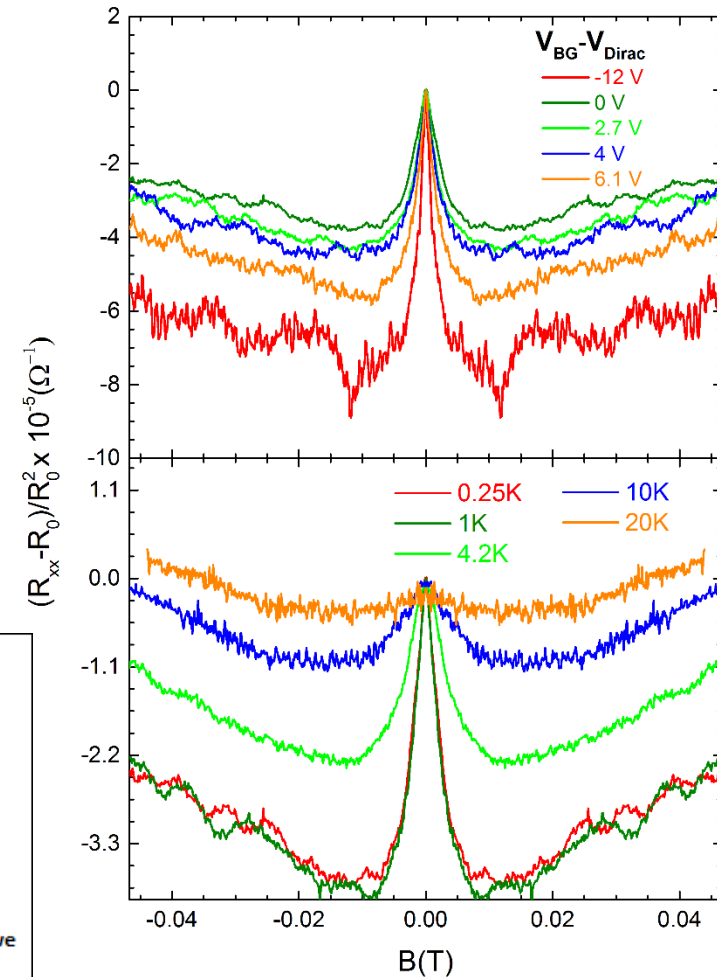
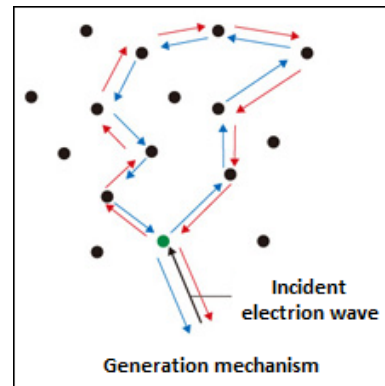
Quantum Hall



S. Xiang, V. Miseikis, L. Planat, S. Guiducci, S. Roddaro, C. Coletti, F. Beltram, S. Heun, Nano research 2016, 9 (6): 1823-1830

Weak localization

- Well established phenomenon due to quantum interference.
- Requires that $L_e < L_\phi$.
- A non-zero magnetic field breaks time reversal symmetry suppressing quantum interference.
- High temperature suppresses quantum interference.



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- E. McCann et al., Phys. Rev. Lett. **97**, 146805 (2006) :

$$\frac{\Delta R_{xx}}{R_0^2} = -\frac{e^2}{\pi h} \left[F\left(\frac{\tau_B^{-1}}{\tau_\phi^{-1}}\right) - F\left(\frac{\tau_B^{-1}}{\tau_\phi^{-1} + \tau_{iv}^{-1}}\right) - 2F\left(\frac{\tau_B^{-1}}{\tau_\phi^{-1} + \tau_*^{-1}}\right) \right]$$

$$\Delta R_{xx} = R_{xx} - R_0$$

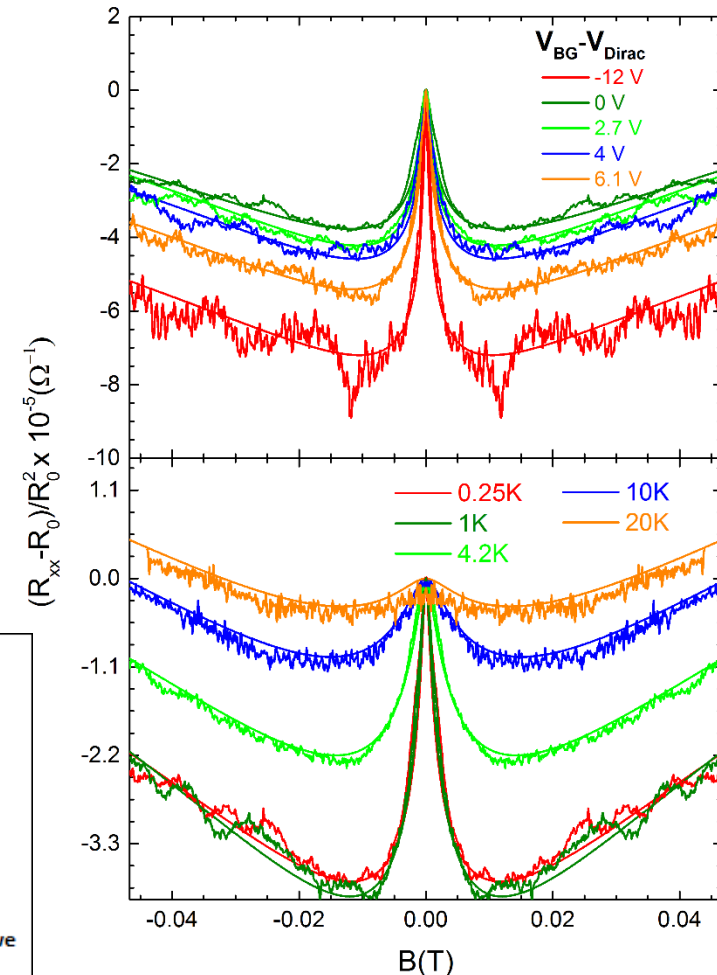
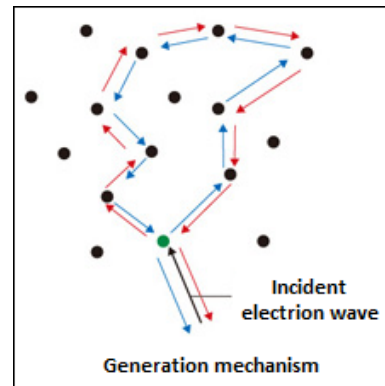
$$\tau_B = 4DeB / \hbar$$

$$F(z) = \ln(z) + \psi(0.5 + z^{-1})$$

τ_ϕ = dephasing time (inelastic)

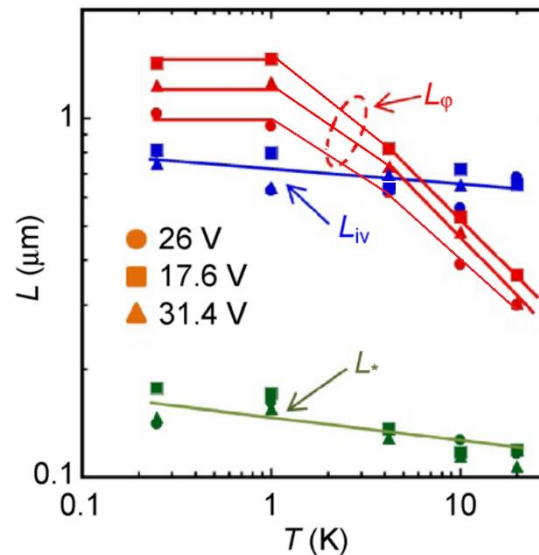
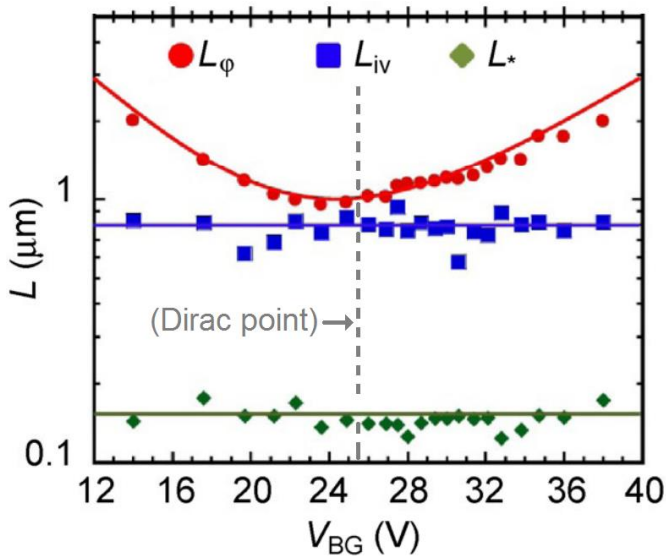
τ_{iv} = intervalley scattering time (elastic)

τ_* = intravalley scattering time (elastic)

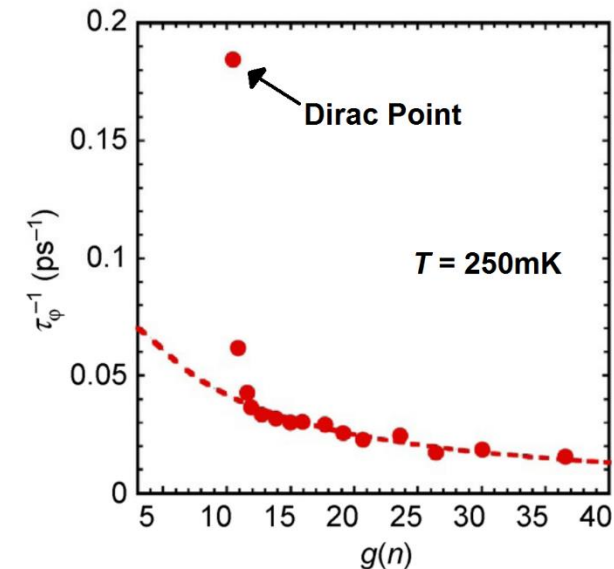
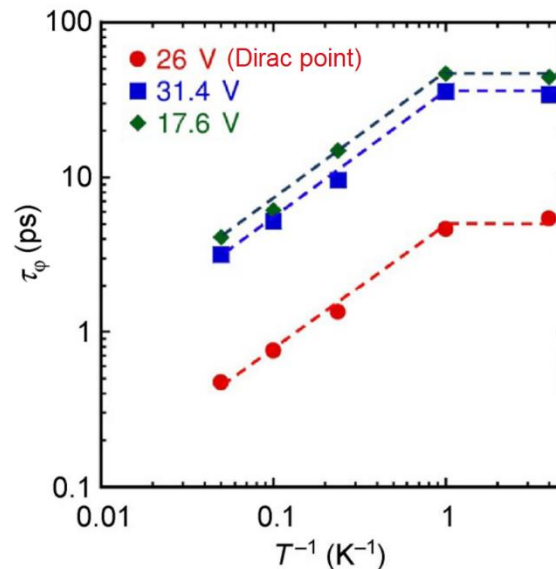
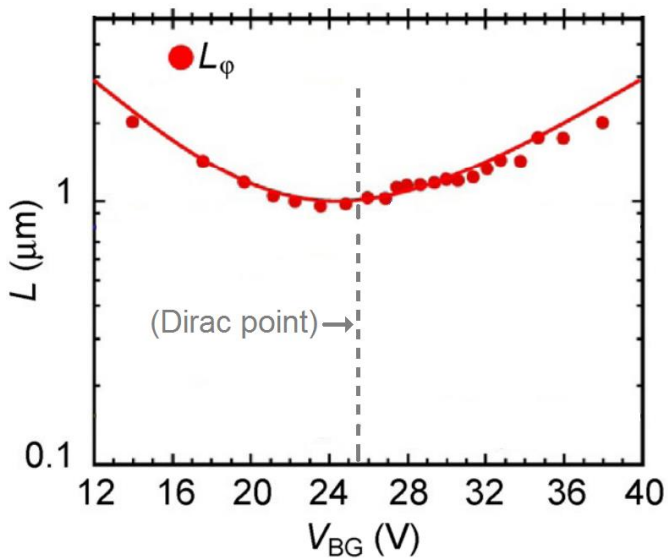


Weak localization

- Both L^* and L_{iv} only weakly depend on V_{BG} and temperature.
- Stronger dependence of L_ϕ from V_{BG} and temperature.
- $L_\phi > 1\mu\text{m}$ (at $T = 0.25\text{K}$).

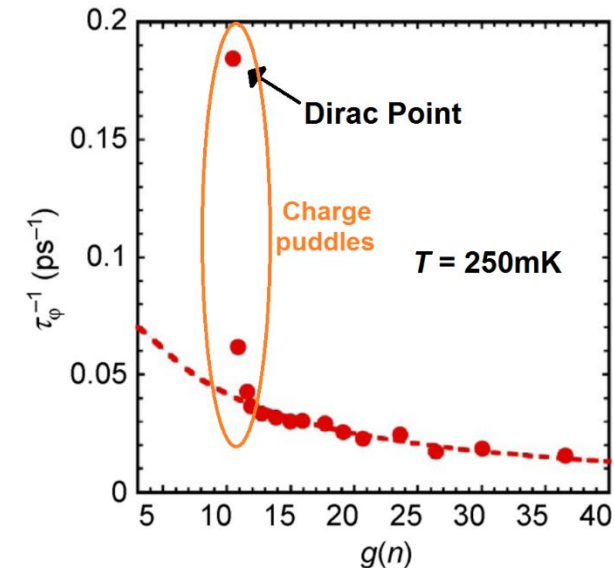
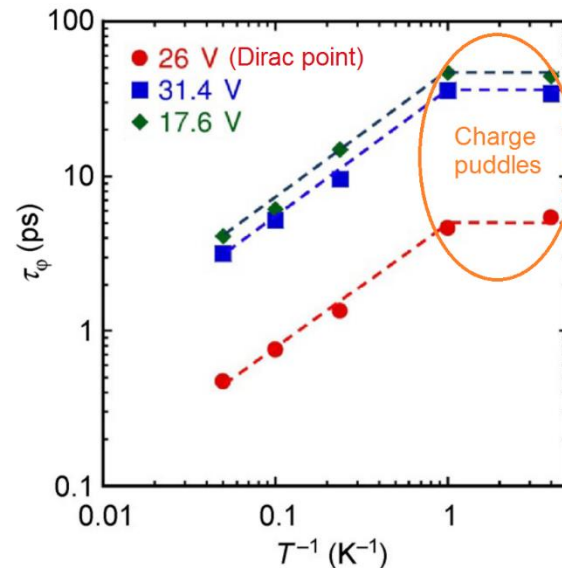
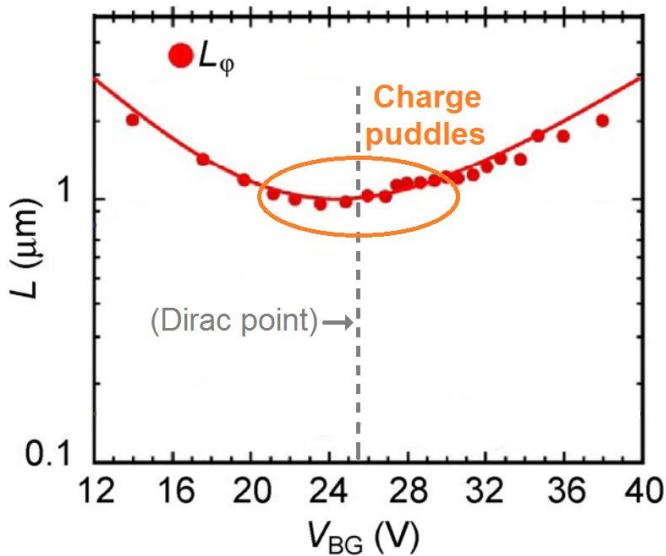


Weak localization



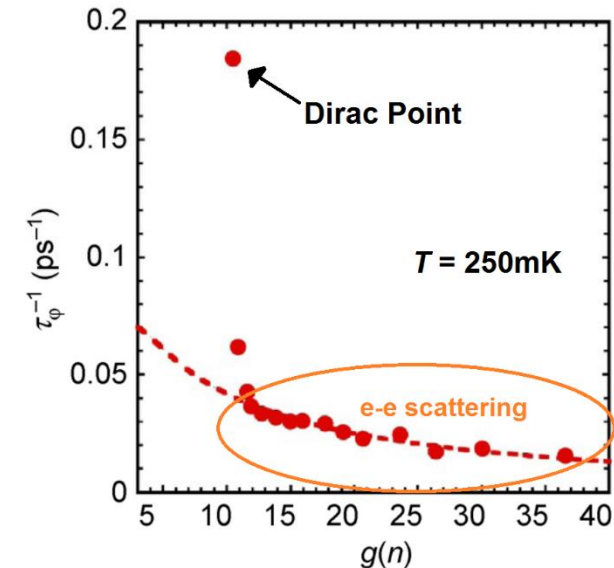
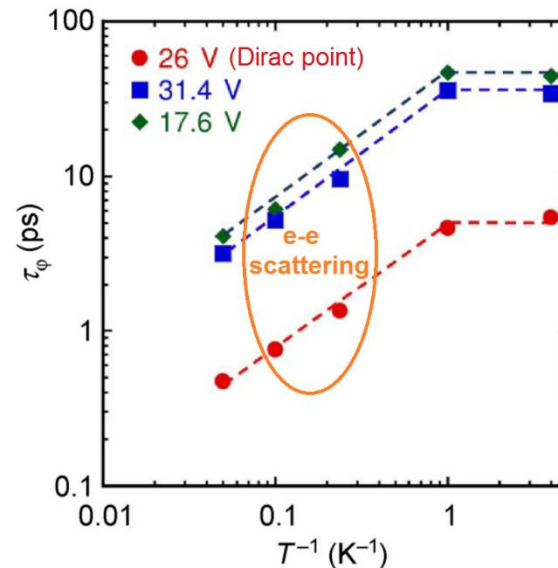
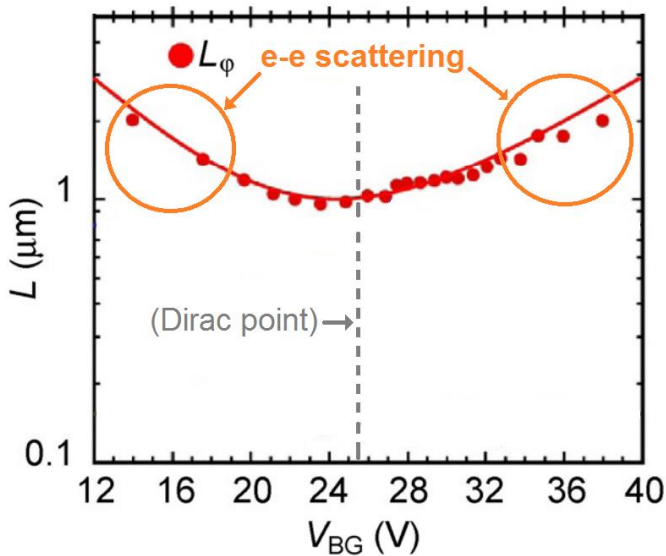
Weak localization

- L_ϕ minimum at the Dirac point (limited by the size of charge puddles).
- Where $\tau_\phi(T^{-1})$ is flat the dephasing time is limited by the charge puddles.
- L_ϕ and τ_ϕ saturate at different values when changing V_{BG} (the size of the charge puddles changes with V_{BG}).
- $\tau_\phi^{-1}(g(n))$ fit with Nyquist formula: where the fit agrees with data e-e is the main inelastic scattering mechanism.



Weak localization

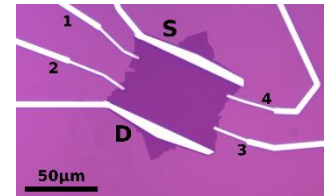
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- $\tau_\phi^{-1}(g(n))$ fit with Nyquist formula: where the fit agrees with data e-e is the main inelastic scattering mechanism.
- In the range where $\tau_\phi \propto T^{-1}$ electron-electron interaction is the main inelastic scattering mechanism.



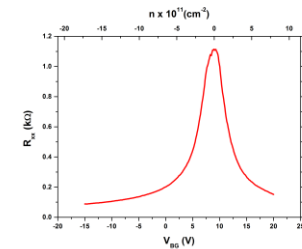
Conclusions



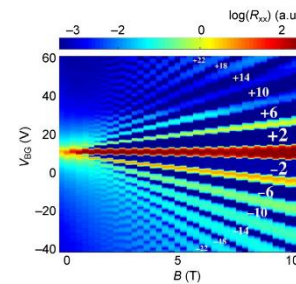
We studied large monocrystalline CVD graphene flakes on SiO₂.



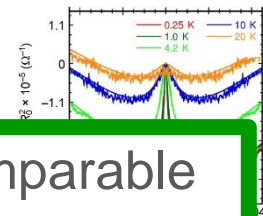
High mobility and low intrinsic charge carriers measured.



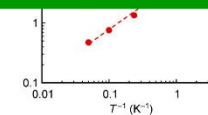
Observed 12 well developed quantum Hall plateaus.



Studied weak localization and measured a dephasing length above 1 μm



The quality of our CVD single crystal graphene is good and comparable to what is measured in exfoliated graphene.





Fundings



Acknowledgements

SGM group

S. Heun



S. Roddaro



S. Xiang



L. Bours



L. Planat

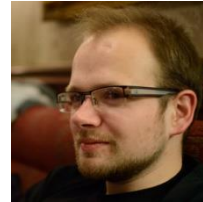


Graphene group (IIT)

C. Coletti



V. Miseikis

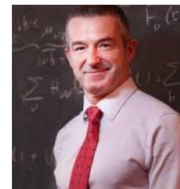


CNR-nano and SNS

L. Sorba



F. Beltram





Thank you for your attention!