Dephasing in Strongly Anisotropic Black Phosphorus

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Catania, 14/12/2016

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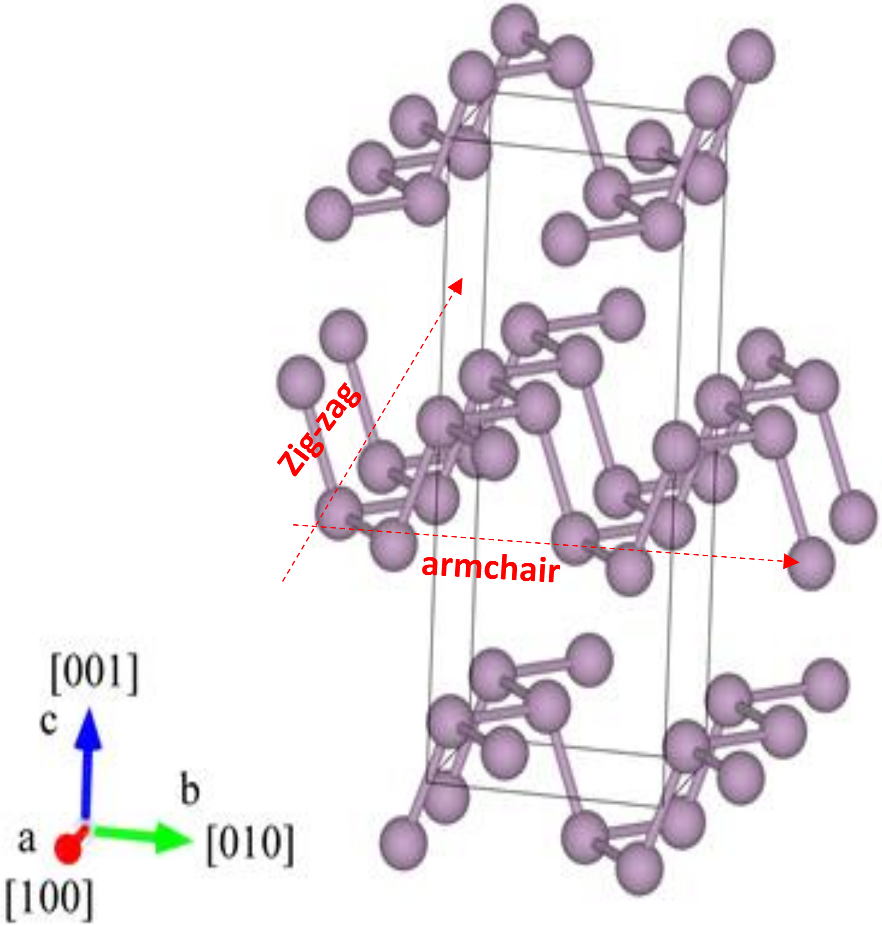
Black phosphorus

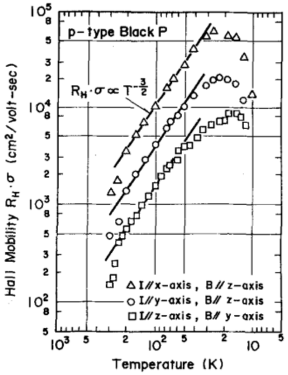


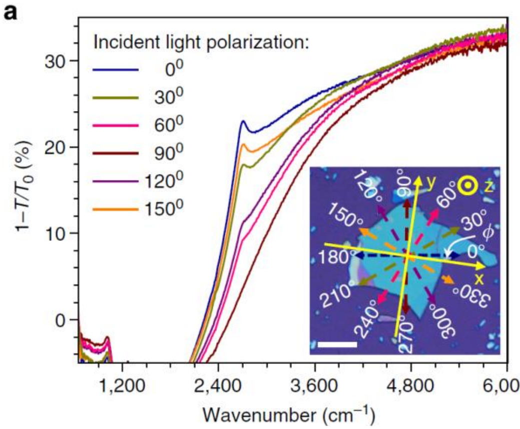
Cell parameters
a=3.13Å
b=10.47Å
c=4.37Å

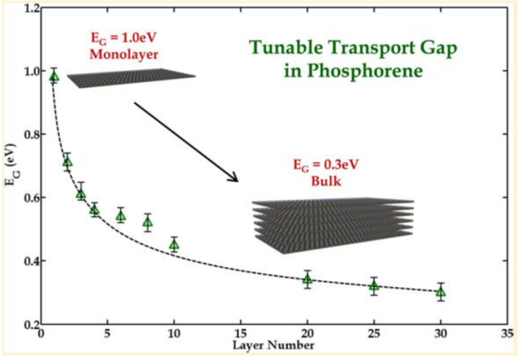
A. Morita, Appl. Phys. A 39 (1986) 227

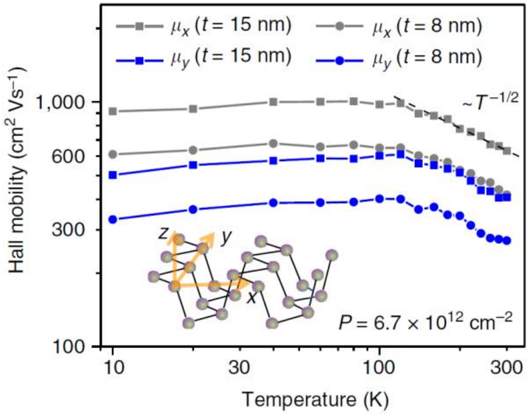
The reinassance of black phosphorus

- ✓ In 1914 first successful synthesis (Bridgman) and in 2007 synthesis at room pressure (Lange, Nilges)
- ✓ p-type semiconductor: 0.3eV direct band gap and high hole mobility (64,000 cm^2/Vs @ 20 K)
- ✓ 1983 (Narita): n-type doping by Te
- ✓ 2014: First publications on bP layered thin films
- ✓ Highly reactive in air
- ✓ Band-gap tunable with layer number
- ✓ **In-plane anisotropy** of optical and transport properties
- ✓ … and much more! SdH, QH…









A. Morita, Appl. Phys. A 39 (1986) 227, S. Das et al., Nano Lett. 14 (2014) 5733, F. Xia et al., Nat. Comm. 5 (2014) 4458

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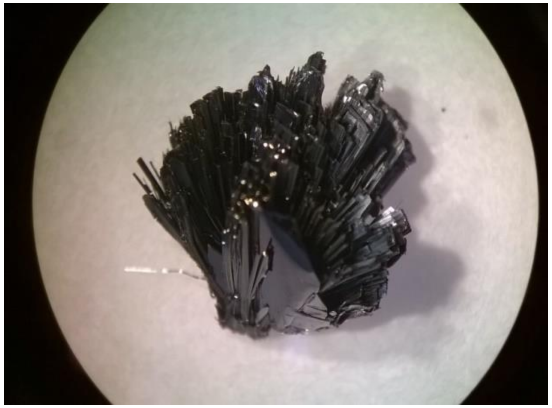
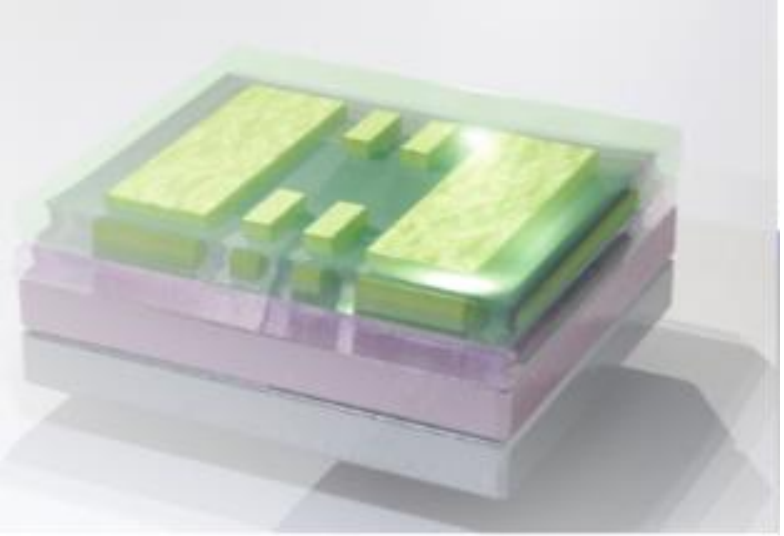
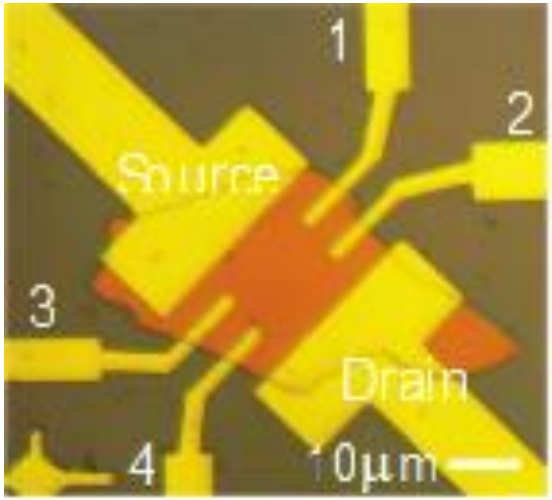
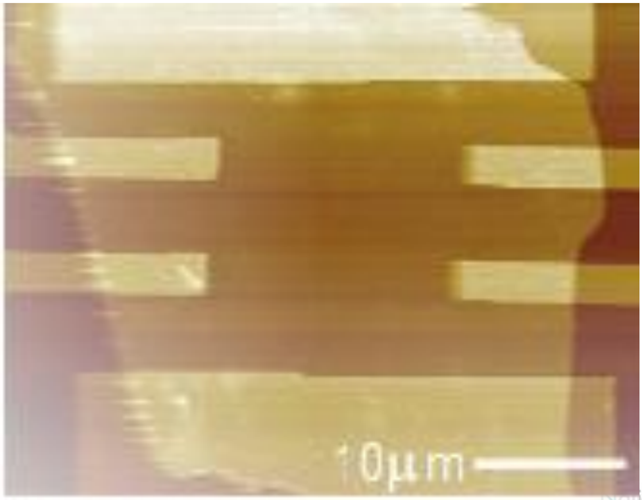
bP Field Effect Transistor





PMMA
MMA
Ti/Au contacts
bP flake
HMDS
SiO_2 thermal oxide
Si



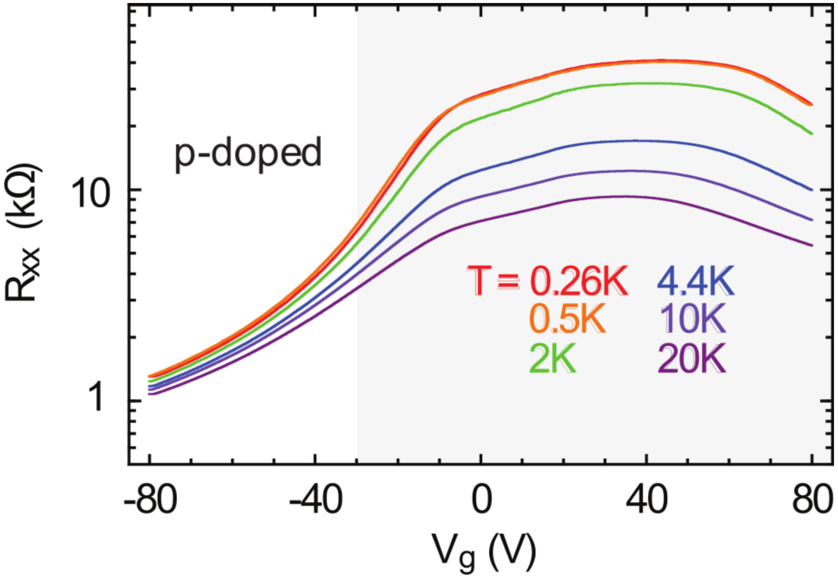


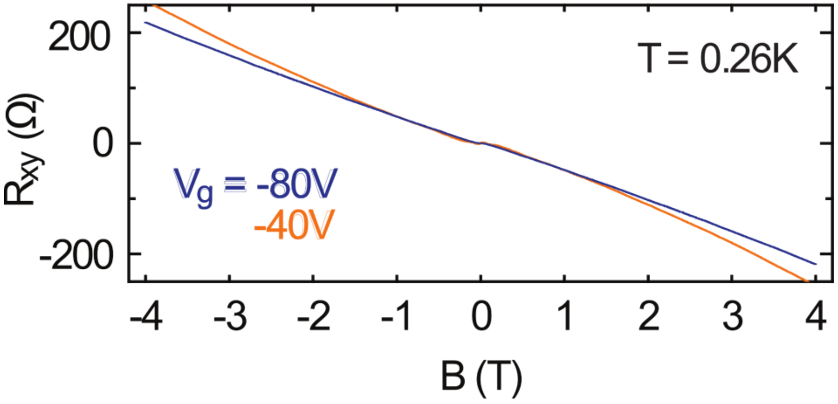
R_{xx}: 1-2
R_{xy}: 1-3

Flake thickness:
65 ± 2 nm

N.Hemsworth et al, PRB 94, 245404, (2016)

Transport Characterization

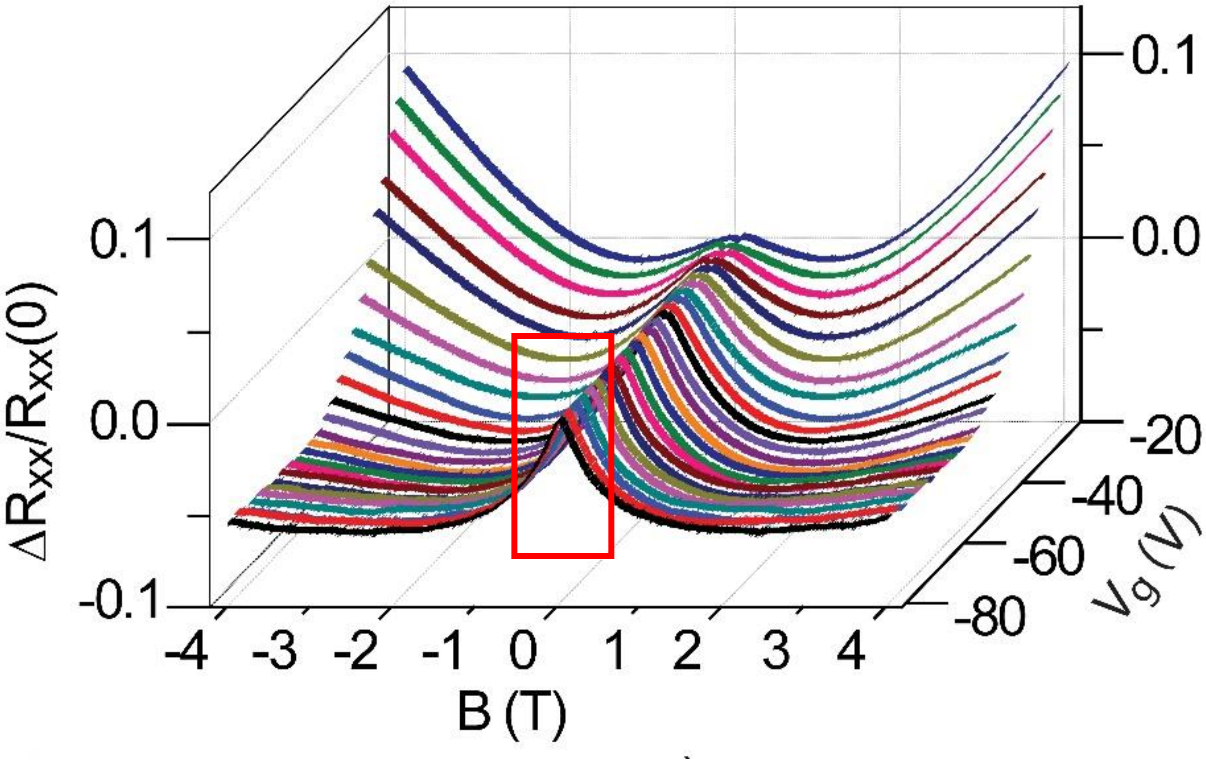




- p ~ V_g for $V_g < -30$ V
- p = 10^{13} cm^{-2} for $V_g = -30$ V
- Field-effect mobility μ: 300 cm^2/Vs at $V_g = -70$ V
- Negligible T-dependence in μ for 0.26 K < T < 20 K

N.Hemsworth et al, PRB 94, 245404, (2016)

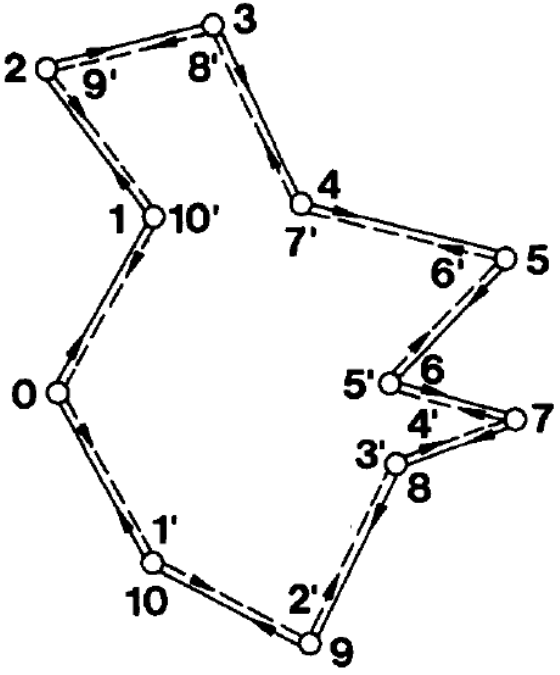
Longitudinal magnetotransport measurements



N.Hemsworth et al, PRB 94, 245404, (2016)

Weak Localization

Weak localization is a quantum effect related to coherent scattering at low temperatures.



Picture from Bergmann, Weak localization in thin films, Physics Reports 107, 1984

Amplitude A_1

Amplitude A_2

Normal Diffusion Model:

$$P = |A_1|^2 + |A_2|^2 = 2\,|A|^2$$

Coherent Addition:

$$P = |A_1 + A_2|^2 = |2\,A|^2 = 4\,|A|^2$$

Since weak localization is a coherent scattering effect:

- It's depressed by magnetic field
- It's smeared by temperature

Weak Localization: Hikami-Larkin-Nagaoka model

$$\Delta\sigma = -\frac{e^2}{2\pi^2\hbar}\left(\Psi\left(\frac{1}{2}+\frac{B_1}{B}\right)-\frac{3}{2}\Psi\left(\frac{1}{2}+\frac{B_2}{B}\right)+\frac{1}{2}\Psi\left(\frac{1}{2}+\frac{B_3}{B}\right)\right)$$

Where Ψ is the digamma function

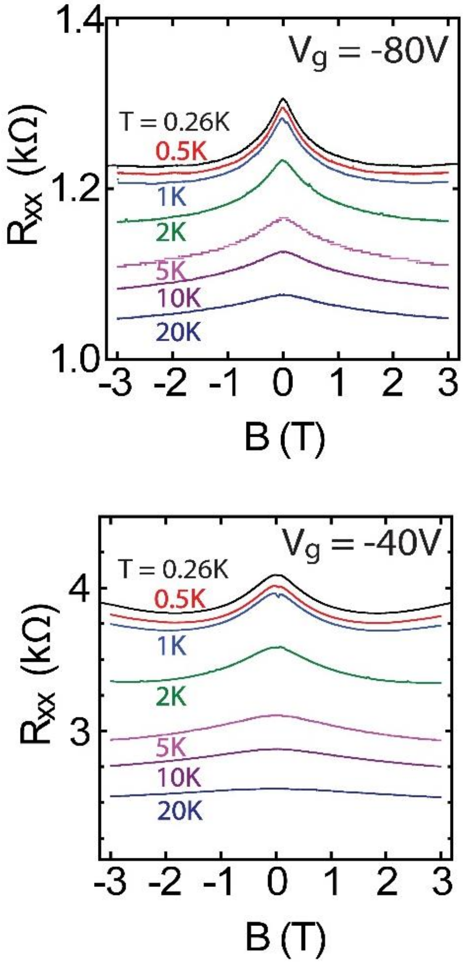
$$B_1 = \boxed{B_0} + \xcancel{B_{so}} + \xcancel{B_s}$$

$$B_2 = \xcancel{\frac{4}{3}}B_{so} + \xcancel{\frac{2}{3}}B_s + \boxed{B_\phi}$$

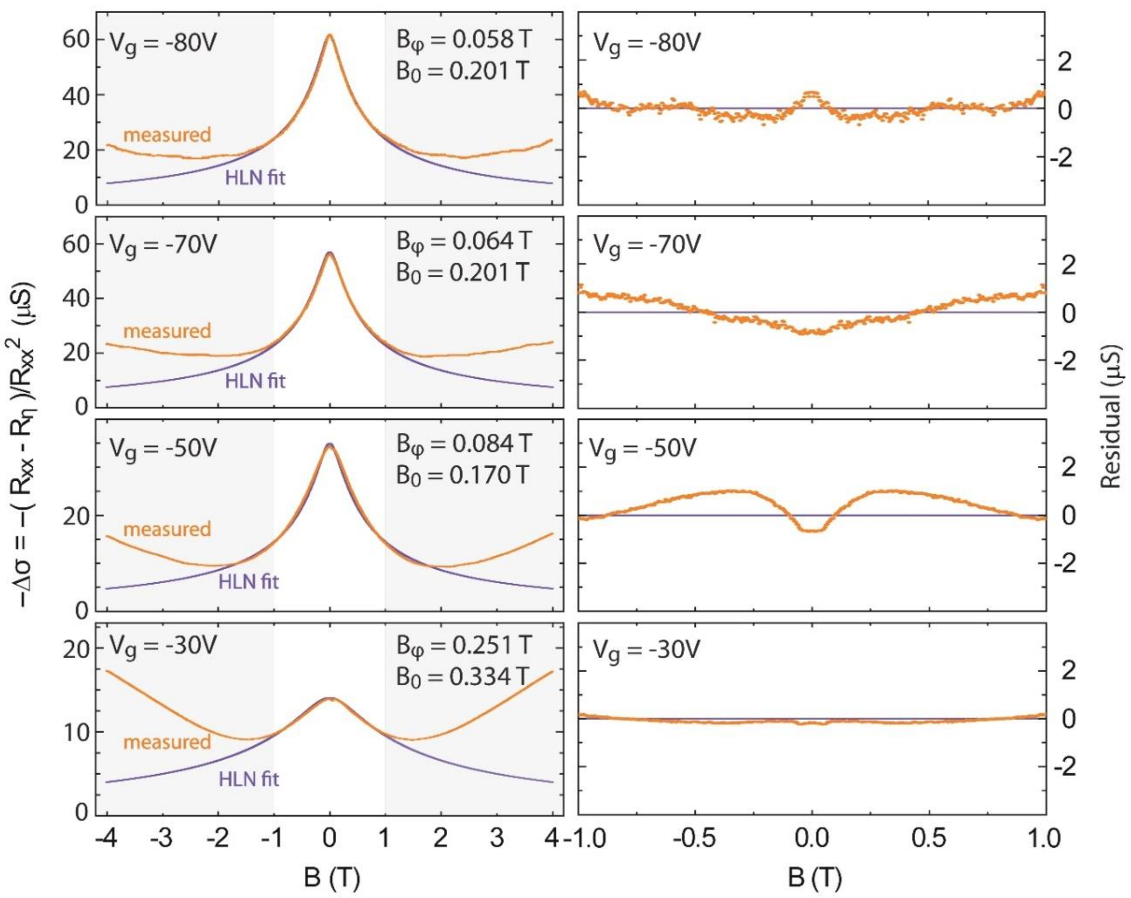
$$B_3 = \xcancel{2}B_s + B_\phi$$

S. Hikami, A. I. Larkin, and Y. Nagaoka,
Prog. Of Theor. Phys. 63 (1980) 707.

Weak Localization



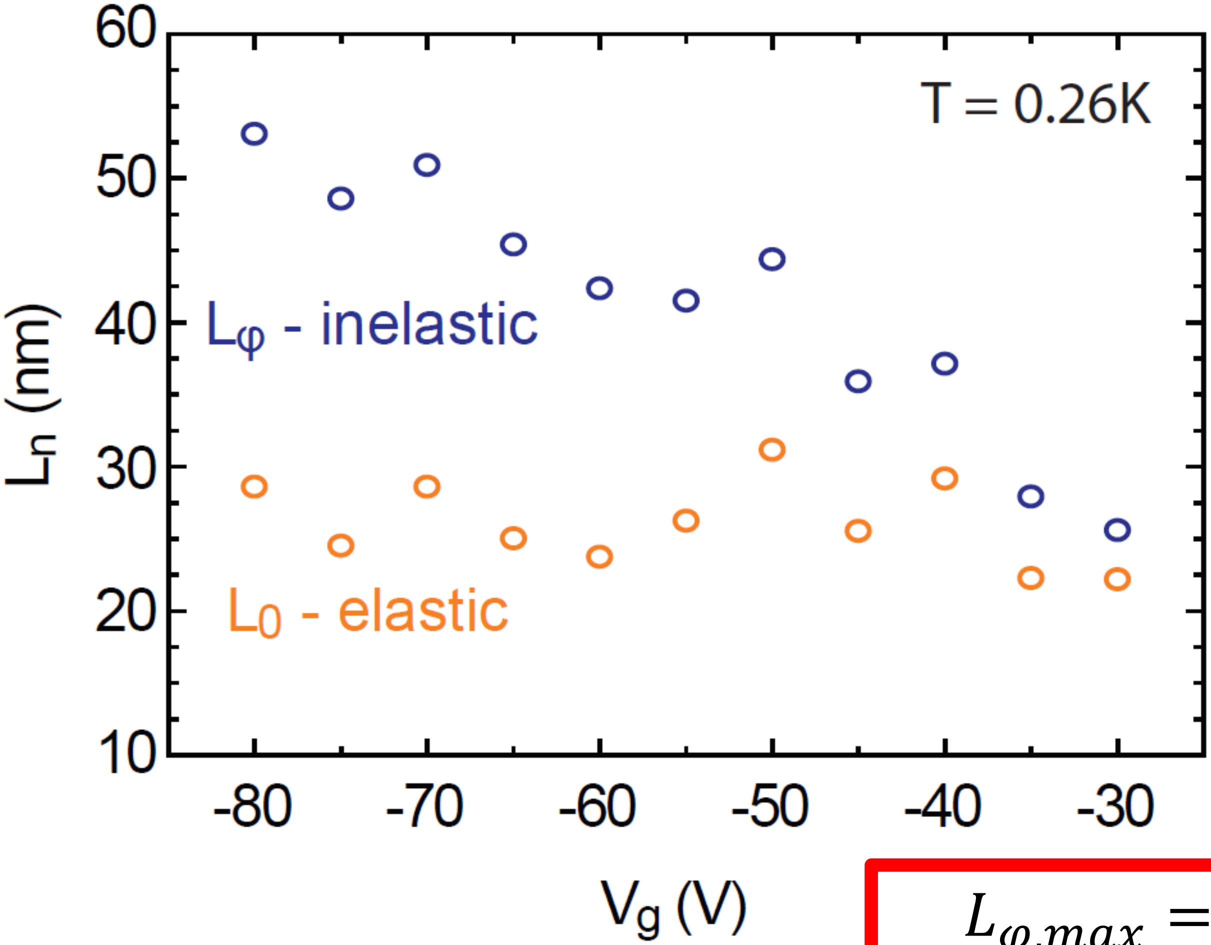
T = 0.26 K



N.Hemsworth et al, PRB 94, 245404, (2016)

Scattering Lengths

$BL^2 = {}^{h}/_{4e}$



$$\boxed{L_{\varphi,max} = 55\ nm}$$

N.Hemsworth et al, PRB 94, 245404, (2016)

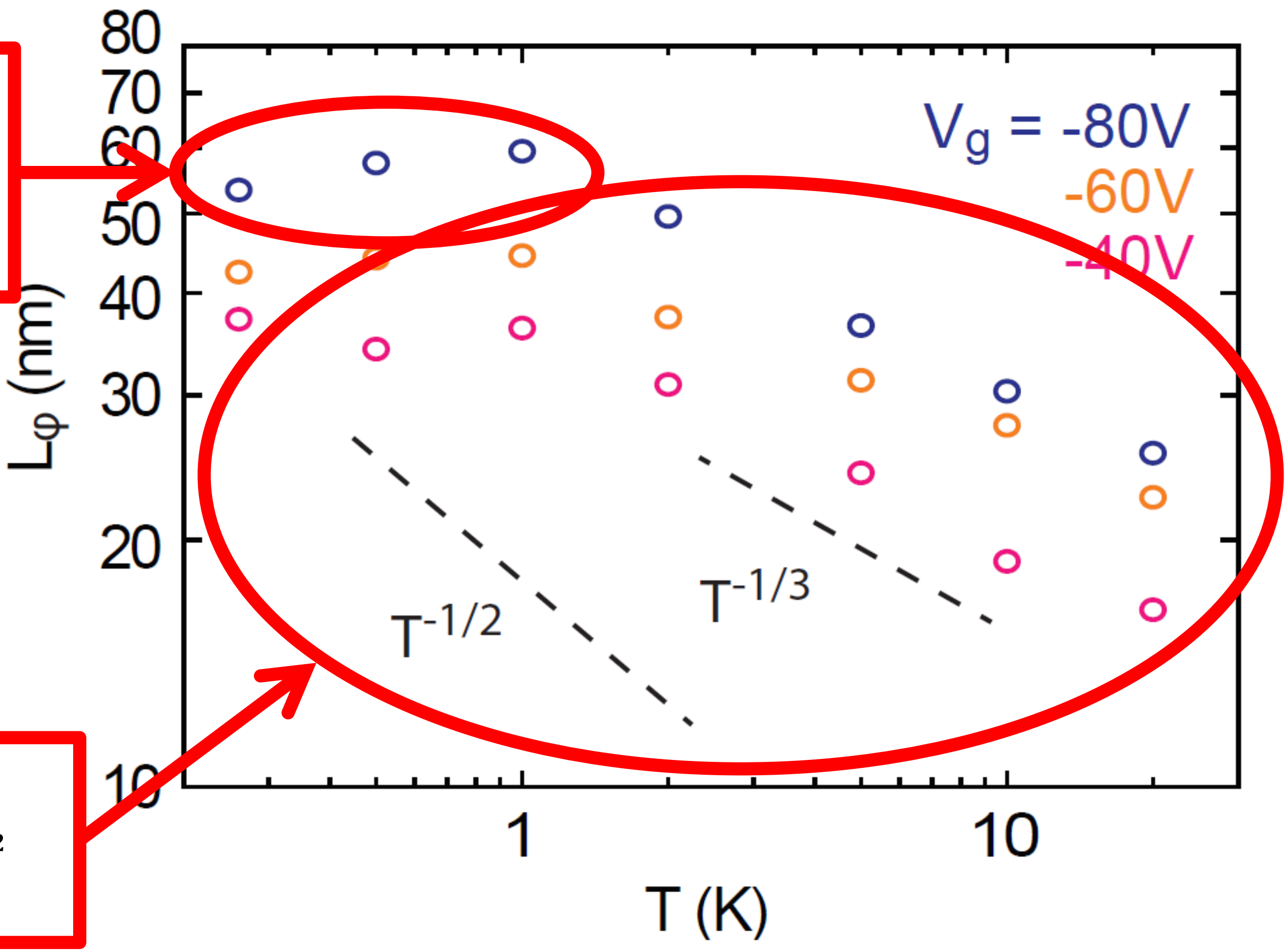
Scattering Lengths

- Dephasing length vs. inelastic scattering time: $L_\varphi = \sqrt{D\tau_\varphi}$ with *D* diffusion coefficient
- Ballistic transport: $\tau_\varphi \propto T^{-2}$ or $L_\varphi \propto T^{-1}$
- Diffusive transport ($L_0 < L_\varphi$):

$$\boxed{\tau_\varphi \propto T^{-1} \text{ or } L_\varphi \propto T^{-1/2}}$$

Lin and Bird, Jour. Phys. Cond. Mat. 14, R501, (2002)

Scattering Lengths

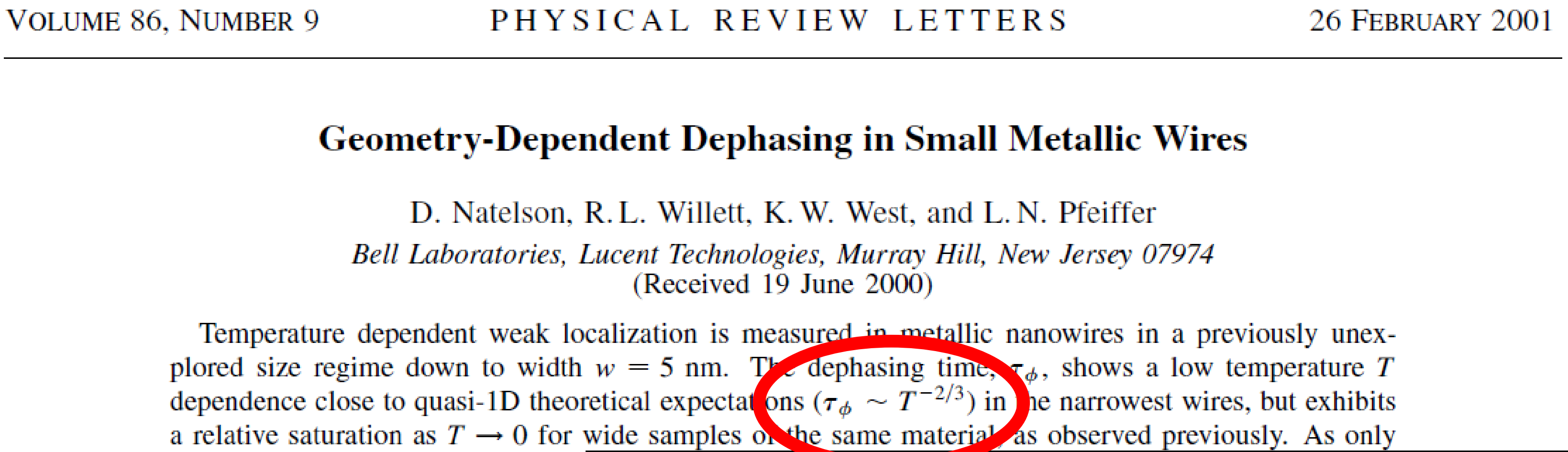


Saturation most likely due to impurities.

L_φ does not follow a $T^{-1/2}$ behaviour.

N.Hemsworth et al, PRB 94, 245404, (2016)

Qasi-1D systems

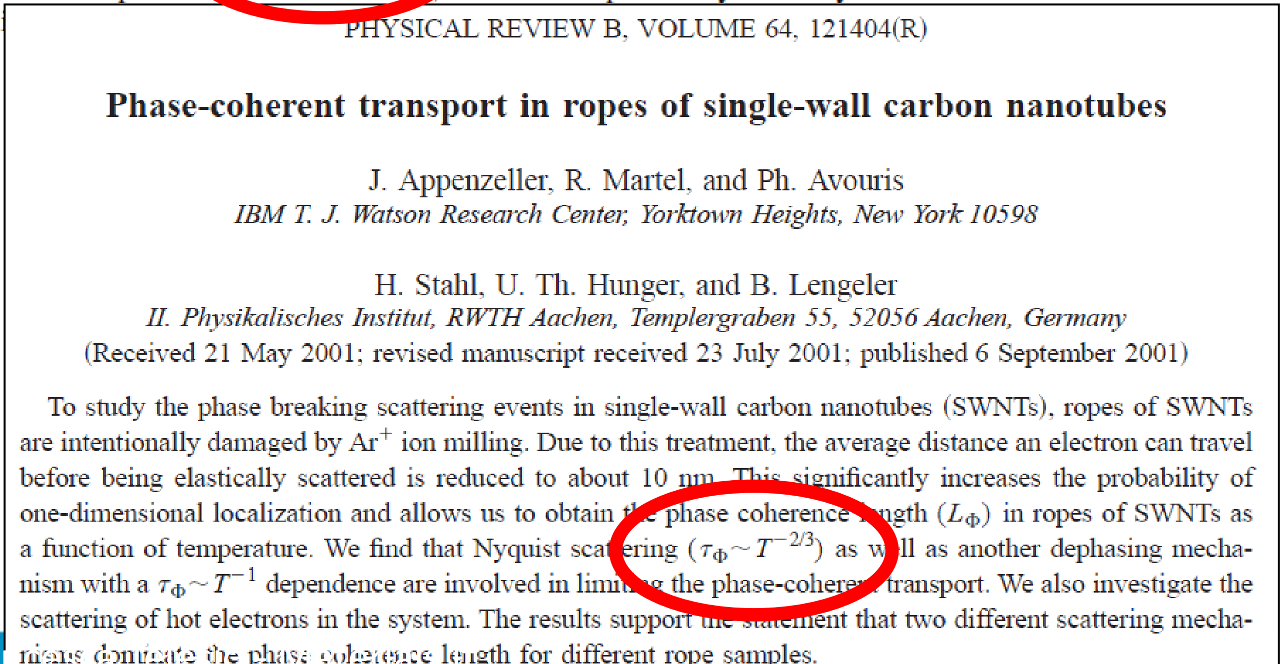


VOLUME 86, NUMBER 9 PHYSICAL REVIEW LETTERS 26 FEBRUARY 2001

Geometry-Dependent Dephasing in Small Metallic Wires

D. Natelson, R. L. Willett, K. W. West, and L. N. Pfeiffer
Bell Laboratories, Lucent Technologies, Murray Hill, New Jersey 07974
(Received 19 June 2000)

Temperature dependent weak localization is measured in metallic nanowires in a previously unexplored size regime down to width $w = 5$ nm. The dephasing time, τ_ϕ, shows a low temperature T dependence close to quasi-1D theoretical expectations ($\tau_\phi \sim T^{-2/3}$) in the narrowest wires, but exhibits a relative saturation as $T \rightarrow 0$ for wide samples of the same material, as observed previously. As only sample geometry is varied to exhib
constraint on models of dephasing



PHYSICAL REVIEW B, VOLUME 64, 121404(R)

Phase-coherent transport in ropes of single-wall carbon nanotubes

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IBM T. J. Watson Research Center, Yorktown Heights, New York 10598

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(Received 21 May 2001; revised manuscript received 23 July 2001; published 6 September 2001)

To study the phase breaking scattering events in single-wall carbon nanotubes (SWNTs), ropes of SWNTs are intentionally damaged by Ar$^+$ ion milling. Due to this treatment, the average distance an electron can travel before being elastically scattered is reduced to about 10 nm. This significantly increases the probability of one-dimensional localization and allows us to obtain the phase coherence length (L_Φ) in ropes of SWNTs as a function of temperature. We find that Nyquist scattering ($\tau_\Phi \sim T^{-2/3}$) as well as another dephasing mechanism with a $\tau_\Phi \sim T^{-1}$ dependence are involved in limiting the phase-coherent transport. We also investigate the scattering of hot electrons in the system. The results support the statement that two different scattering mechanisms dominate the phase coherence length for different rope samples.

Comparison with quasi-1D wires

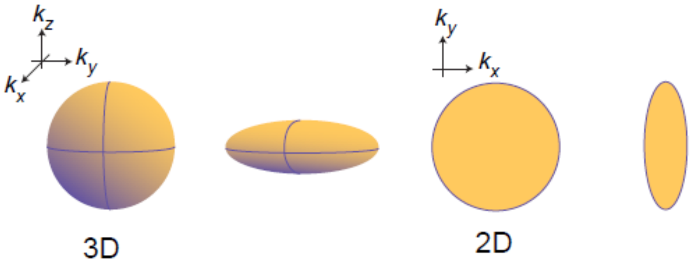
D. Natelson et al.
PRL 86 (2009):

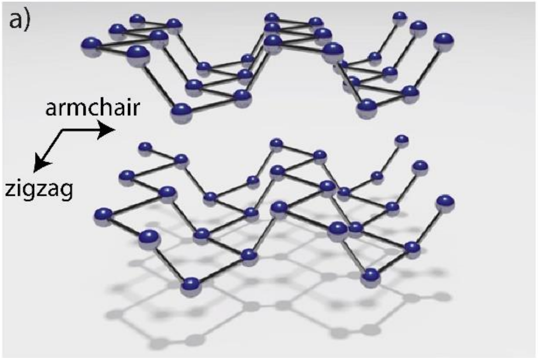
quasi-1D:
$L_\varphi, L_T > w, t$

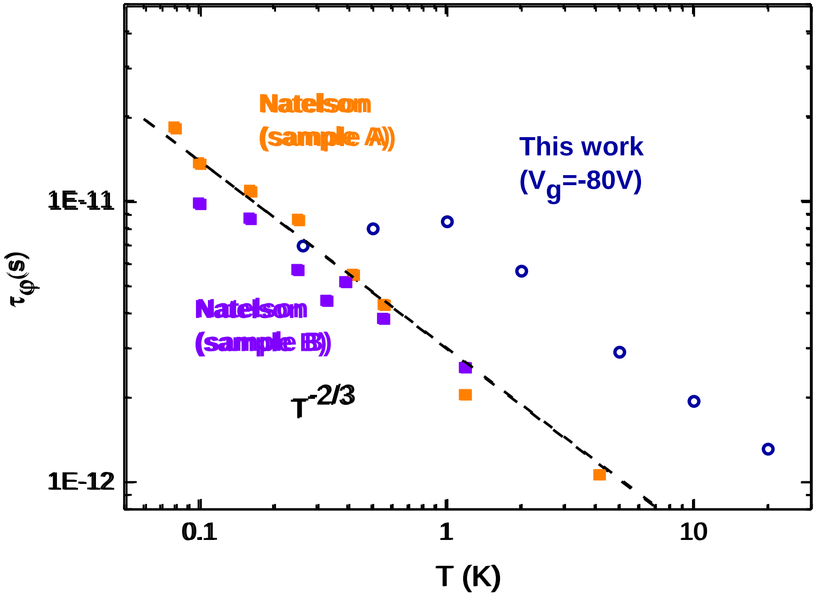
width w
thickness t

$\tau_\varphi \propto T^{-2/3}$

$L_\varphi \propto T^{-1/3}$







$L_\varphi = 55\ nm$

thermal length:

$L_T = \sqrt{\hbar D / k_B T}$
= 10 – 60 nm

N.Hemsworth et al, PRB 94, 245404, (2016)

Conclusions

- ✓Weak localization oberved in a bP FET
- ✓Excellent agreement with HLN model
- ✓Dephasing length L_{φ} reaches 55 nm
- ✓T-dependence of L_{φ} close to quasi-1D
- ✓This seems a further proof of strong in plane anisotropy of bP

N.Hemsworth, et al, PRB 94, 245404, (2016)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*"**Phos**phorene **fun**ctionalization: a new platform for advanced multifunctional materials"*

Thank you for your attention!

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